

**The Gyrokinetic Regime**

**Geometry**

**Velocity Space**

**Linear How-To**

**Units**

**Non-Linear Issues**

**Miscellany**

# Coordinates

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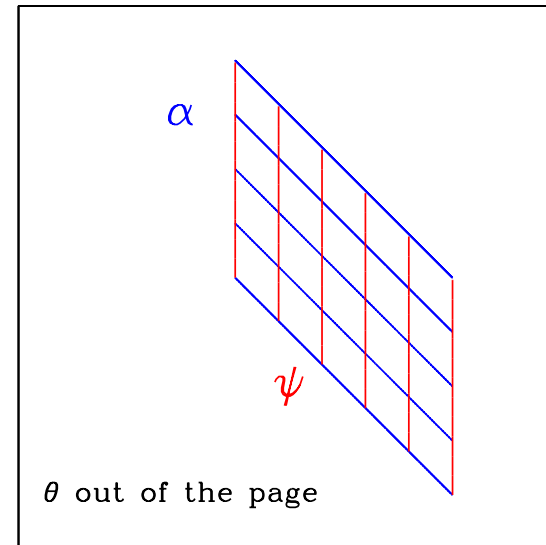
- Given a magnetic field in Clebsch representation

$$\mathbf{B} = \nabla\alpha \times \nabla\psi$$

- Natural perpendicular coordinates are  $\psi$  and  $\alpha$ ; distance along the field line is  $\theta$

- In general,  $(\psi, \alpha, \theta)$  is non-orthogonal

- Nor are they automatically single-valued.



# Simplest Toroidal Limit

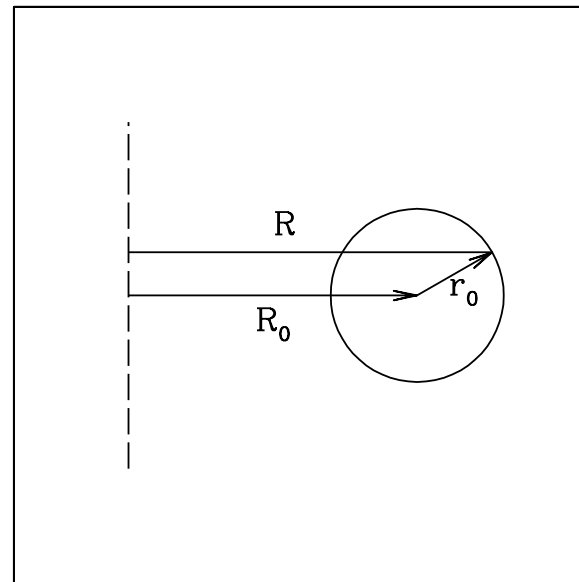
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- Concentric circular flux surfaces:

$$\alpha = \phi - q\left(\theta - \frac{r}{R_0} \sin \theta\right)$$

$$\psi = \psi_0 + (r - r_0) \frac{B_0 r_0}{q}$$

- Note  $\theta$  is the distance along the field line.



(Concentric circles)

# Flux Tube Limit

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- Simulation domain is some number of gyroradii in  $\psi$  and  $\alpha$  directions, and  $\sim 2\pi qR$  along field line

- Take advantage of ordering to write

$$h = \hat{h}(\theta) \exp(iS)$$

(Fluctuations vary slowly along field line, rapidly across)

- Make contact with ballooning theory, defining

$$S = n_0 [\alpha + q(\Psi)\theta_0]$$

where  $n_0$  is the integer toroidal mode number.

# Theta Dependencies of $\nabla B$ Drift

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- Work out  $\theta$  dependencies in various terms
- $\nabla B$  drift:

$$\frac{v_{\perp}^2}{2} \frac{1}{\Omega B^2} \mathbf{B} \times \nabla B \cdot \nabla S = \left( \frac{k_{\theta} \rho_i}{2} \right) \frac{v_{\perp}^2}{2} \left[ \omega_{\nabla B} + \omega_{\nabla B}^{(0)} \theta_0 \right],$$

where

$$\omega_{\nabla B} = \frac{2}{B^2} \frac{d\Psi}{d\rho} \hat{\mathbf{b}} \times \nabla B \cdot \nabla \alpha, \quad \omega_{\nabla B}^{(0)} = \frac{2}{B^2} \frac{d\Psi}{d\rho} \hat{\mathbf{b}} \times \nabla B \cdot \nabla q.$$

- In the *high aspect ratio* concentric circles limit:

$$\omega_{\nabla B} = \frac{2a}{R_0} (\cos \theta + \hat{s} \theta \sin \theta) \quad \omega_{\nabla B}^{(0)} = -\frac{2a}{R_0} (\hat{s} \sin \theta)$$

# Definitions

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- To get the previous expressions, define the normalized flux surface coordinate  $\rho = r/a$  and also  $k_\theta$ :

$$k_\theta \equiv \frac{n_0}{a} \frac{d\rho}{d\Psi_N} = \frac{n_0 q_0}{r_0}$$

- Normalized poloidal flux  $\Psi_N \equiv \Psi / (a^2 B_a)$
- Here,  $a$  is half the diameter of the LCFS at the elevation of the magnetic axis.
- The field  $B_a$  is normalized to the toroidal field at  $R_a$ , where  $R_a$  is the avg of the min and max R on the LCFS.

# Theta Dependencies of Curvature Drift

- Curvature drift:

$$\omega_{\kappa} = \omega_{\nabla B} + \frac{8\pi}{B^2} \frac{dp}{d\rho}.$$

where  $B = B(\theta)$  and  $p$  is the total pressure.

- In the *high aspect ratio* concentric circles limit  $\omega_{\nabla B} = \omega_{\kappa}$
- In general, the curvature drift is always *bad* on the outboard midplane of a tokamak, but the  $\nabla B$  drift can be reversed.

## $\theta$ Dependencies of Parallel Derivatives

- Parallel derivatives:

$$\hat{\mathbf{b}} \cdot \nabla = \frac{\mathbf{B}_0 \cdot \nabla}{B_0}$$

- In the *high aspect ratio* concentric circles limit  $\hat{\mathbf{b}} \cdot \nabla = a/(qR_0)$
- Freedom in the definition of  $\theta$  can be exploited to remove the  $\theta$  dependence from this operator (equal\_arc = T)



## $\theta$ Dependencies of Perp Gradients

- Bessel functions have argument  $\gamma = |\nabla S|v_{\perp}/\Omega$ ; always enters as square ( $\gamma^2$ ).
- General expression:

$$\begin{aligned} |\nabla S|^2 &= \frac{n_0^2}{a^2} |\nabla (\alpha + q\theta_0)|^2 \\ &= k_{\theta}^2 \left( \frac{d\Psi}{d\rho} \right)^2 \left| \nabla \alpha \cdot \nabla \alpha + 2\theta_0 \nabla \alpha \cdot \nabla q + \theta_0^2 \nabla q \cdot \nabla q \right|. \end{aligned}$$

- In the *high aspect ratio* concentric circles limit

$$|\nabla S|^2 = k_{\theta}^2 \left| 1 + \hat{s}^2 \theta^2 - 2\theta\theta_0 \hat{s}^2 + \theta_0^2 \hat{s}^2 \right|$$

# Flux-Surface Averages

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- Given Jacobian  $J$ , flux-surface average of a quantity  $\Gamma$  is

$$\langle \Gamma \rangle = \lim_{\Delta\rho \rightarrow 0} \frac{\int \Gamma J d\theta d\alpha d\rho}{\int J d\theta d\alpha d\rho}.$$

- Flux-surface avg of a radially directed quantity (a flux) is

$$Q_{\text{sim}} = \frac{\langle \mathbf{Q} \cdot \nabla \rho \rangle}{\langle \nabla \rho \rangle}$$

which will appear in the transport equation as

$$\frac{3}{2} \frac{d}{dt} \langle nT \rangle + \frac{1}{V'} \frac{d}{d\rho} A Q_{\text{sim}} + \dots = 0,$$

where the surface area  $A = 2\pi \langle |\nabla \rho| \rangle \int J d\theta$ .

# Loose Ends

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- Note the limit in this expression:

$$\langle \Gamma \rangle = \lim_{\Delta\rho \rightarrow 0} \frac{\int \Gamma J d\theta d\alpha d\rho}{\int J d\theta d\alpha d\rho}.$$

- In a flux-tube simulation, since the perpendicular box size is measured in gyroradii and  $\rho_* \ll 1$ , it is appropriate to take the integral over the entire simulation domain.
- In other words, a flux-tube simulation naturally calculates the flux-surface average of the various fluxes.

# Critical Elements for Axisymmetric Users

- Tokamak user has three choices:
  1. Assume shifted circles & specify  $r/a$ ,  $a/R$ ,  $q$ ,  $\hat{s}$ , and  $\alpha$
  2. Specify the shape of the flux surface and  $B_p(\theta)$
  3. Read in numerical specification from MHD equilibrium solver
- In each case, the user can use the two free functions of the Grad-Shafranov equation to vary the magnetic shear and the pressure gradient freely: details controlled by **bishop**

## Shifted Circles (theta\_grid\_parameters)

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- Magnetic shear:  $\text{shat} = \hat{s}$
- Aspect ratio:  $\text{eps1} = 2a/R_0$
- Safety factor:  $\text{pk} = 2a/(qR_0)$ ; *i.e.*,  $q = \text{eps1}/\text{pk}$
- Minor radius:  $\text{eps} = r/R$
- Shafranov shift:  $\text{shift} = \alpha = -2Rq^2 d\beta/dr$   
(not the coordinate  $\alpha$ !)

# Local Equilibrium

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- Surface of constant  $\Psi$  specified by:

$$R_N(\theta) = R_{0N}(\rho) + \rho \cos [\theta + \delta(\rho) \sin \theta],$$

$$Z_N(\theta) = \kappa(\rho)\rho \sin \theta.$$

(More general expressions easily implemented.)

- Poloidal field determined from

$$B_p(\theta) = \frac{|\nabla\Psi|}{R} = \frac{d\Psi}{d\rho} \frac{|\nabla\rho|}{R}$$

- Conventions:  $R_N = R/a$ , etc.,

$$R_{0N}(\rho) = R_{0N}(\rho_c) + R'_{0N} d\rho, \quad \delta(\rho) = \delta(\rho_c) + \delta' d\rho, \quad \kappa(\rho) = \kappa(\rho_c) + \kappa' d\rho$$

# Local Equilibrium: Details I

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- Input  $q$ , note that  $\oint \alpha d\theta = -2\pi q$  and then determine  $d\Psi/d\rho$  from

$$\frac{d\Psi}{d\rho} = \frac{I}{2\pi q} \oint \frac{d\theta}{R^2} (\nabla\theta \times \nabla\rho \cdot \nabla\phi)^{-1} = \frac{I}{2\pi q} \oint \frac{J d\theta}{R^2}$$

- Define three integrals:

$$A(\theta) = \int J d\theta \left[ \frac{1}{R^2} + \left( \frac{I}{B_p R^2} \right)^2 \right], \quad B(\theta) = I \int J d\theta \left[ \frac{1}{(B_p R)^2} \right],$$

$$C(\theta) = I \int J d\theta \left[ \frac{\sin u + R/R_c}{B_p R^4} \right]$$

## Local Equilibrium: Definitions

- $u(\theta)$ : angle between the horizontal and the tangent to the magnetic surface in the poloidal plane
- $R_c(\theta)$ : local radius of curvature of surface in the poloidal plane.
- From papers by Mercier and Luc, used later by C. Bishop and R. Miller



## Local Equilibrium: Details II

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- Upon defining  $\bar{A} = \oint d\theta \dots$ , one can show

$$\hat{s} = \frac{\rho dq}{q d\rho} = \frac{\rho}{2\pi q} \frac{d\Psi}{d\rho} (\bar{A}I' + \bar{B}p' + 2\bar{C})$$

- Primes are derivatives w.r.t.  $\Psi$
- User thus may specify any two of  $p'$ ,  $I'$ , and  $\hat{s}$ . Several input options determined by [bishop](#)

## Local Eq Inputs (theta\_grid\_parameters)

- Magnetic shear:  $\text{shat} = \hat{s}$  (for bishop = 1)
- Minor radius:  $\text{rhoc} = \rho_c$  (Normalized by a!)
- Safety factor:  $\text{qinp} = q$
- Elongation:  $\text{akappa} = \kappa$
- Deriv of elongation:  $\text{akappri} = d\kappa/d\rho$

## Local Eq Inputs (theta\_grid\_parameters)

- Triangularity:  $\text{tri} = \delta$
- Deriv of triangularity:  $\text{tripri} = d\delta/d\rho$
- Center of flux surf:  $\text{Rmaj} = R_{0N}$
- Local shift:  $\text{shift} = dR_{0N}/d\rho < 0$
- Center of LCFS:  $\text{R\_geo} = R_{\text{geo}N}$  (for normalization only)

## Eq Inputs (theta\_grid\_eik\_knobs)

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- These parameters may be set whether the equilibrium information comes from a local model *or* from numerical data (a file)
- Magnetic shear:  $s\_hat\_input = \hat{s}$   
(for `bishop > 1`)
- Pressure gradient:  $beta\_prime\_input = d\beta/d\rho < 0$   
(for `bishop = 4`)

## Eq Selections (theta\_grid\_eik\_knobs)

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- Set at most one to true:  
`ppl_eq`, `transp_eq`, `efit_eq`, `gen_eq`, `vmom_eq`, `local_eq`
- Set `iflux = 1` to use numerical equilibrium, `0` otherwise
- Choose definition of  $\rho$  (flux surface label):  
`irho = 1` ... ( $\rho = \sqrt{\Phi/\Phi_a}$ )  
`irho = 2` ... ( $\rho = d/D$ )  
`irho = 3` ... ( $\rho = \Psi/\Psi_a$ )
- Likely never change: `itor = 1`

## Eq Recommendations (theta\_grid\_eik\_knobs)

- Watch out for the units, esp normalizations by  $a$
- Set `writelots = T` (it's essentially free information)
- Set `isym = 0` (allow up-down asymmetry)
- Set `equal_arc = F` (and avoid a rare bug)
- Set `del_rho = 1.e-3` (nearly always adequate)

## Numerical Equilibria (theta\_grid\_eik\_knobs)

- If given the choice, use inverse solutions  $R(\Psi, \theta)$ ,  $Z(\Psi, \theta)$  for accuracy
- Specify file with data using `eqfile = '...'`  
No blanks allowed, but other characters okay.
- Don't forget to set `bishop`. Use `1` to get the data from the file, `> 1` to alter it.

# References

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- M. D. Kruskal and R. M. Kulsrud, Phys. Fluids, 1:265, 1958
- M. A. Beer and S. C. Cowley and G. W. Hammett, Phys. Plasmas, 2:2687, 1995
- R. L. Miller, *et al.*, Phys. Plasmas, 5:973, 1998
- Notes for this talk: [http://gk.umd.edu/g\\_short.pdf](http://gk.umd.edu/g_short.pdf)